

## § 4.2 Compactifications of 6d (1,0) SCFT's

The  $\mathbb{Z}_k$  orbifold of  $\mathcal{N}=2$  linear quivers

Want to look at compactifications of

$T_k^N(1,0)$  SCFT's:  $N$  M5 branes probing  
 $A_{k-1}$  singularity  
 "A-type conformal matter"

brane setup:

M-th.:

	$\mathbb{C}^2/\mathbb{Z}_k$										
	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$	$x^{10}$
NM5	$x$	$x$	$x$	$x$	$x$						$x$

↓ identify  $x^{10}$  with M-theory circle

IIA:

	$\mathbb{C}^2/\mathbb{Z}_k$									
	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$
ND4	$x$	$x$	$x$	$x$	$x$					

→  $(\mathcal{N}=2$  5d  $U(N)$  SYM) /  $\mathbb{Z}_k$

scalars:  $x^9, x^5 + ix^6 = \gamma_1, x^7 + ix^8 = \gamma_2$

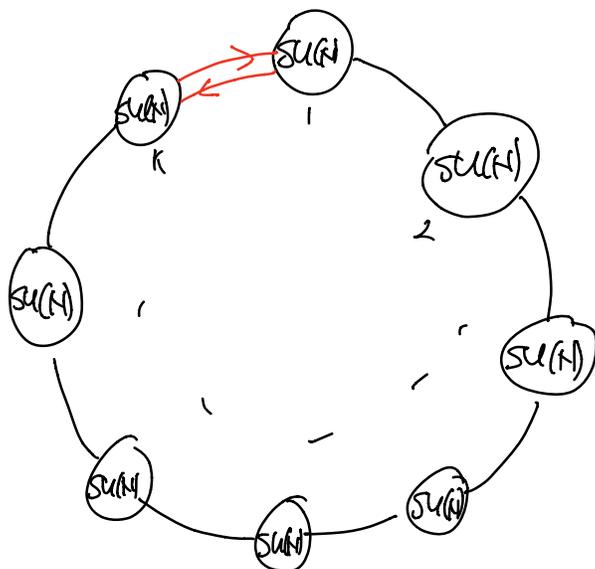
$\mathbb{Z}_k$ :







resulting 5d quiver theory:



5d theory has  $U(1)^{2k}$  global symmetry

$$\subset SU(k)_S \times SU(k)_Y \times U(1)_T$$

Non-abelian symmetry is broken by Wilson lines for  $SU(k)_S \times SU(k)_Y$

→ perturbative gauge theory description in 5d

$$\tau = \sum_{i=1}^k \tau_i = \sum_{i=1}^k \frac{4\pi^2}{g_{YM,i}^2} = \frac{1}{R}$$

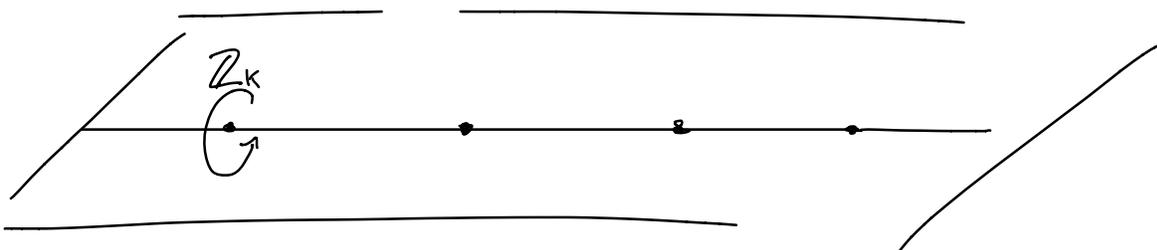
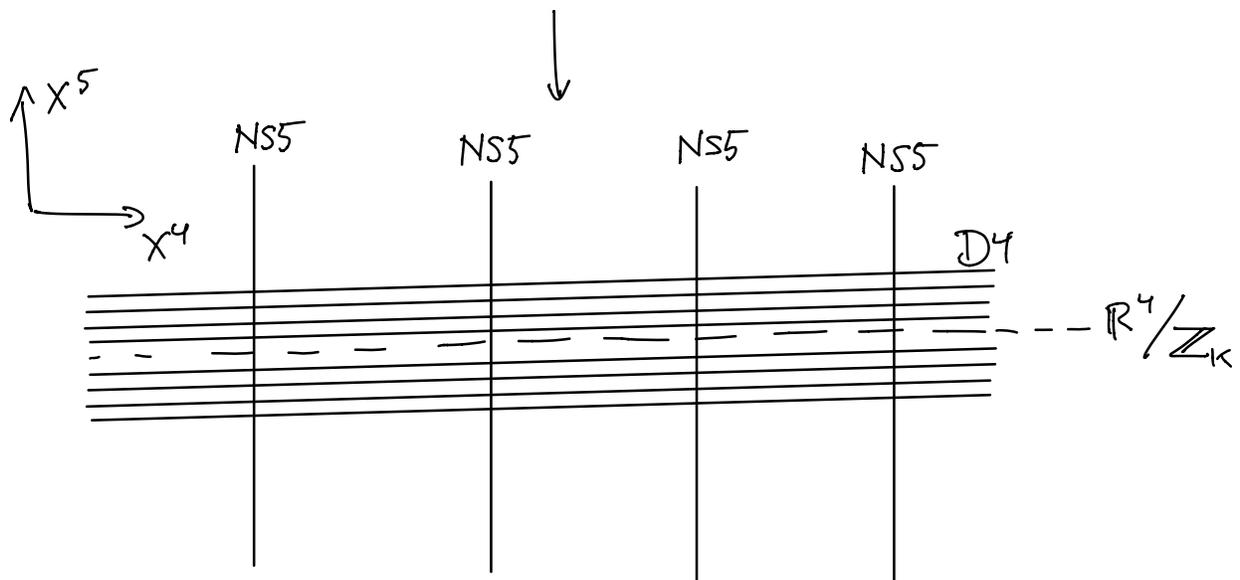
$\tau_i = \int_{C_i} \mathcal{B}$ , where  $C_i$  are 2-cycles in the resolution of  $A_{k-1}$ -sing.   
 ↖ B-field   
 ↗ radius of  $X^{10}$ -circle

Further compactification on  $S^1$  would give

4d  $\mathcal{N}=2$  theory

→ to get  $\mathcal{N}=1$  gauge th. in 4d,  
need to include NS5-branes:

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$\mathbb{C}^2/\mathbb{Z}_k$				$x^9$
ND4	x	x	x	x	x					
n NS5	x	x	x	x		x	x			

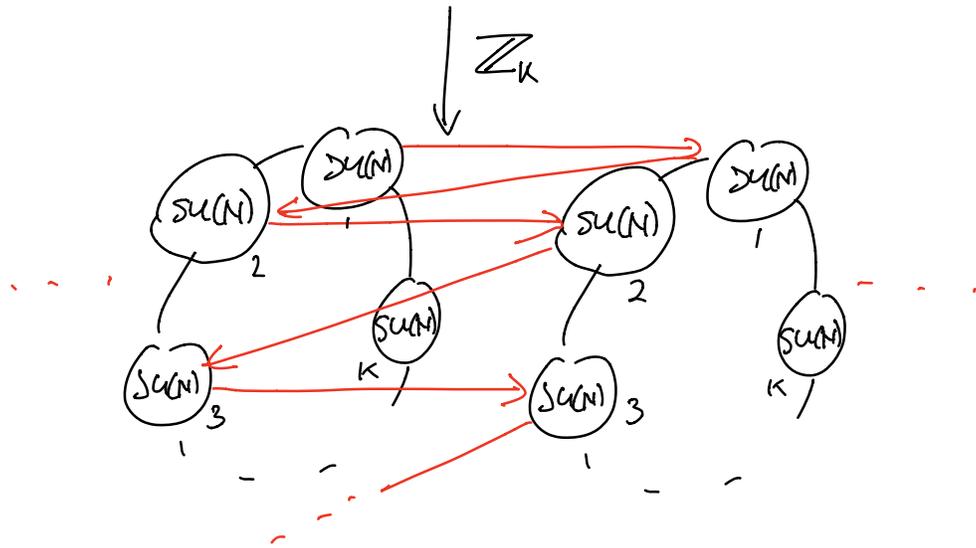
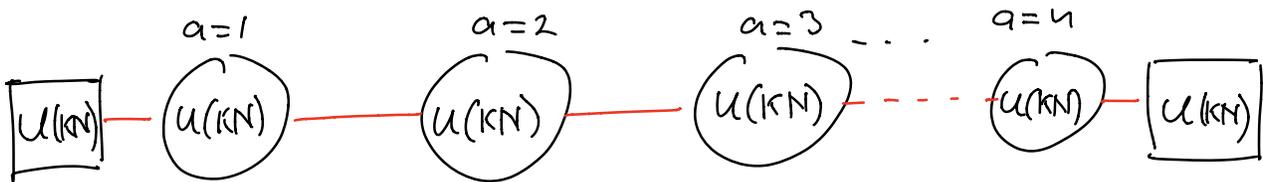


→  $kN$  fractional D4-branes intersected by  $n$  NS5-branes

→ ( $\mathcal{N}=2$  linear quiver of 4d  $U(kN)$  gauge groups)  $/\mathbb{Z}_k$

↓  
 $\mathcal{N}=1$  SYSY

by embedding  $\mathbb{Z}_k$  into  $SU(2)_R \times U(1)$  R-symmetry

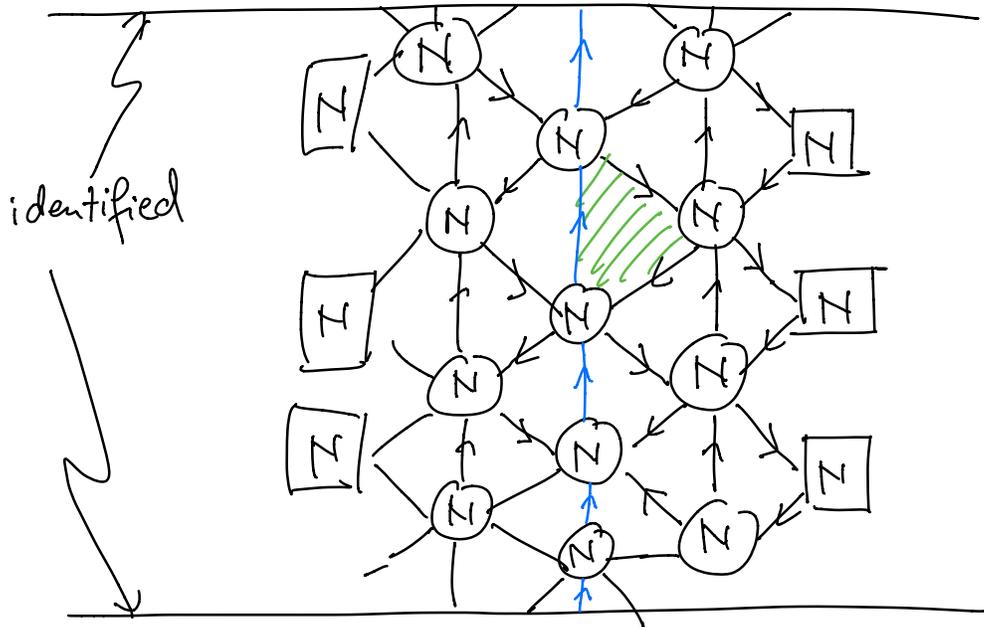


necklace  $\mathcal{N}_a$   $\mathcal{N}_{a+1}$

each  $\mathcal{N}=2$  hyper-multiplets splits to 2  
 $\mathcal{N}=1$  chiral multiplets going from

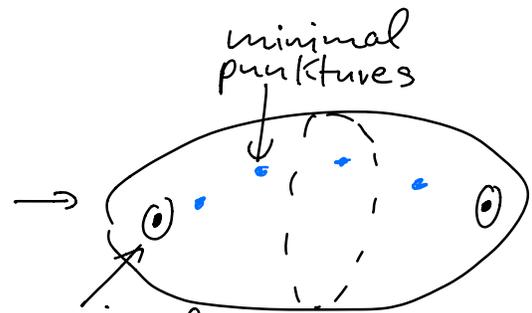
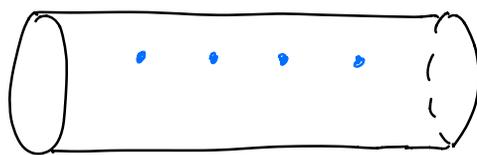
$\mathcal{N}_{a,i} \rightarrow \mathcal{N}_{a+1,i}$  and  $\mathcal{N}_{a+1,i} \rightarrow \mathcal{N}_{a,i+1}$

Resulting structure is tessellation of cylinder



triangular faces associated to cubic superpotentials

6d-lift of 4d theory:



maximal punctures with  $SU(N)^k$  flavor-sym.

each minimal puncture carries  $U(1)_\alpha$  global symmetry

→  $n+1$   $U(1)$ 's

Altogether :  $U(1)^{2k+n}$  global symmetries

remaining  $U(1)^{2k-1} \leftarrow$  Cartan generators  
of  $SU(k)_L \times SU(k)_S \times U(1)_T$

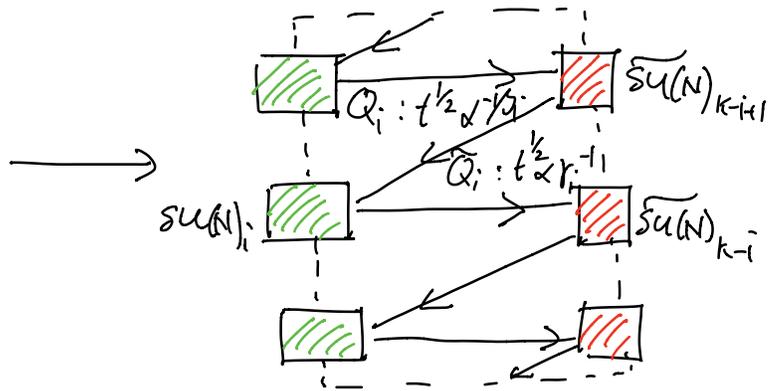
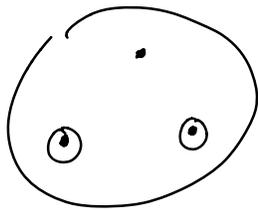
$$U(1)_T \times \left[ \frac{U(1)^k}{U(1)} \right]_S \times \left[ \frac{U(1)^k}{U(1)} \right]_R$$

$\rightarrow$   $n$  marginal couplings:

relative positions of  $n+1$  minimal punctures

The free trinion

sphere with 3 punctures



$2kN^2$  free chiral fields :

- $Q_{b_i}^{a_i}$  in  $SU(N)_i \times \overline{SU(N)}_{k-i+1}$

- and  $\tilde{Q}_{a_i}^{b_{i+1}}$  in  $\overline{SU(N)}_i \times SU(N)_{k-i}$

$k$   $SU(N)_i$  and  $k$   $\overline{SU(N)}_i$  global symmetries

$\rightarrow$  maximal punctures